

**EXPERIMENTAL INVESTIGATION OF NEUTRON PRODUCTION
AND TRANSMUTATION OF ^{129}I , ^{237}Np , ^{238}Pu , ^{239}Pu ON PB-TARGET
IN “ENERGY PLUS TRANSMUTATION” SETUP IRRADIATED BY
2.52 GEV DEUTERONS IN NUCLOTRON (JINR, DUBNA)**

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ABSTRACT

The experiment is a part of the scientific program „Investigation of physical aspects of electronuclear method of energy production and transmutation of radioactive waste using relativistic beams from the JINR Synchrophasotron/Nuclotron“ – under the name of project „Energy plus Transmutation“ (Journal Kerntechnik, 2003, V.68, p.p. 48-55). Results of the first experiment with deuteron beam at the energy of 2.52 GeV are given in this paper. Samples of isotopes of ^{129}I ,

²³⁷Np, ²³⁸Pu and ²³⁹Pu are gathered in notable amounts in nuclear reactors. There are also produced in setups of industries which use nuclear materials and nuclear technologies. The samples were irradiated in the field of neutrons produced in a lead target and propagated in the uranium blanket. The estimation of its transmutation (radioecological aspect) was obtained in result of measurements of their gamma activities. The information about space-energy distribution of neutrons in the volume of the lead target (diameter 8.4 cm, length 45.6 cm) and the uranium blanket (weight of 206.4 kg natural uranium) was obtained with help of activation threshold detectors (Al, Co, Y, I, Au, Bi and other), solid state nuclear track detectors, He-3 neutron detectors and nuclear emulsions. Comparison of the experimental data with the results of simulation with the MCNPX program was performed.

Key words: Spallation sources, transmutation

Introduction

The project is called “Energy plus Transmutation” (see Krivopustov et al. [1-4]). Scientific description of the project, including main ideas, history, performed experiments’ description and results, uranium calorimeter description, experimental methodology used for neutron and proton field properties investigation (activation and solid state nuclear track detectors, nuclear emulsions, He-3 detectors, threshold detectors, etc.), could be found in the publications of the “Energy plus Transmutation” collaboration [3-8] and in the overview of JINR research by Baldin, Malakhov and Syssakian [9a]. During 1999-2004 various experiments were held with “Energy plus Transmutation” assembly with proton beams in the range of energies from 0.7 GeV to 2.0 GeV. The experiments were focused on general aspects of energy generation by future Accelerator Driven Systems (ADS), e. g., neutron generation and multiplication, neutron spectra determination, generation of secondary isotopes inside Pb- target and U-blanket, energy generation and deposition, neutron induced transmutation of long-lived minor-actinides (²³⁷Np, ²⁴¹Am), fission products (¹²⁹I), and plutonium isotopes (^{238,239}Pu) [3-17]. These investigations appear to be very important for development of ADS usable for future nuclear energy utilization and nuclear fuel cycle safety. This technology has recently attracted considerable attention [18-25]. The use of the deuteron beam was motivated by the possibility of comparison of the data of neutron generation in our set-up with data of Tolstov [24], who used Pb-slot of 50x50x80 cm³, and Vassilkov [25], who used the cylinder with diameters from 16 to 20 cm and length from 60 to 76 cm. This paper describes the experiment with deuteron beam with energy of 2.52 GeV, which was held in JINR Dubna, on 30th November 2005 using superconductivity accelerator Nuclotron by Vexler and Baldin Laboratory High Energies (see Kovalenko et al. [9b])

Experimental setup

General scheme of “Energy plus Transmutation” facility [2-6], which was built in 1998-1999 for spending fuel isotopes is given in figs 1 and 2. The detailed technical design was carried out by the All-Russian Institute of Nuclear Energy Machine Building (VNIAM) in Moscow and manufacturing of the steel structure was performed at the mechanical workshop of the LHE JINR. The “Energy plus Transmutation” setup consists of the following system:

- Lead target divided into four sections (diameter of 84 mm and length of 456 mm, weight of 28.6 kg).

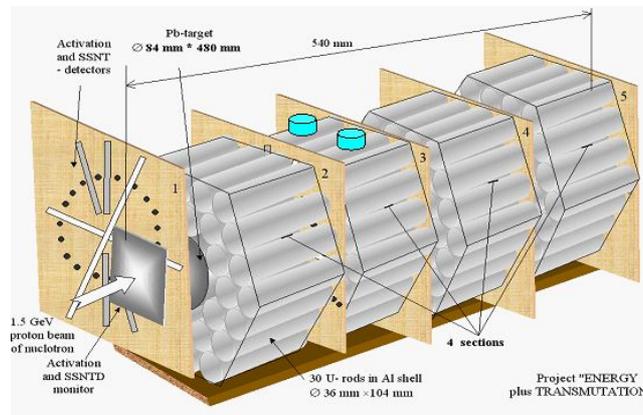


Fig. 1. Scheme of the four-section “Energy plus Transmutation” setup with a massive lead target and uranium blanket [3-5]

- Uranium blanket also divided into four sections; each section consists of 30 fuel rods of natural uranium inside the aluminium cover (34 mm diameter, 104 mm length, weight of 1.72 kg). Each section contains 51.6 kg of uranium, so the whole blanket contains 206.4 kg of natural uranium.
- Beam monitoring system of activation and solid state detectors and proportional ionization chambers.
- He-3 detector system. This kind of detectors was used to determine spatial and energy distributions of neutron fluence [26].
- Set of radioactive samples for transmutation studies. This set contains ^{129}I , ^{237}Np , ^{238}Pu , and ^{239}Pu . Each isotope was hermetically packed inside the duralumin container .
- Backplate for radioactive samples and other foil-based detectors fixation to the top of the 2nd section of the uranium blanket (fig. 1).
- Five plates for activation detectors and Solid State Nuclear Track Detectors (SSNTD) fixation made from special polyethylene-foil.
- Five spectrometers based on nuclear emulsions for neutron registration by proton recoils [5].

- Set of thermometers (thermocouples, thermoresistors, etc.) for determination of the heat generation inside the uranium blanket [3].
- Shielding box made from granulated polyethylene with boron carbide, with a cadmium cover and the outside box made from wood. Box has the dimensions of 100x106x111 cm and weight of 950 kg and can be moved to the irradiation place (focus F3N of the Nuclotron experimental complex) using special rail system. System of activation and threshold detectors, nuclear emulsion, SSNTD, He-3 detectors, and thermal detector system [3, 5, 13] are in general called uranium fission calorimeter [3].

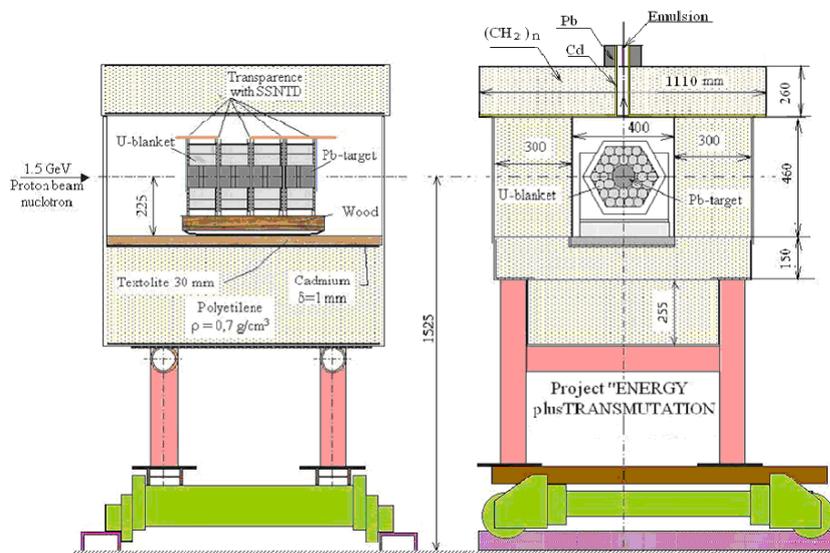


Fig 2. Technical details of the U/Pb-assembly inside a massive shielding and placed into a mobile platform, which can be moved into and out of the beam line. The left side of this figure gives a cross section of the assembly along the deuteron beam line, the right-side shows a cut through the assembly perpendicular to the deuteron beam line in the position [3-5].

Transmutation samples

The transmutation samples (^{129}I , ^{237}Np , ^{238}Pu and ^{239}Pu) were placed on the top of the 2nd section of the uranium blanket (fig.1) fixed on the special paper backplate ($104 \times 140 \times 1 \text{ mm}^3$). In each experiment only one sample of each isotope was used plus one sample with ^{127}I , which was irradiated to subtract its effect in ^{129}I sample, which contains 15% of ^{127}I . Also the ^{238}Pu sample contains some other plutonium isotopes, mainly ^{239}Pu (16.75 %). Radioactive materials are covered by aluminium (special duralumin alloy) with diameter of 34 mm. Some properties of the samples used in 2.52 GeV deuterons are given in table 1.

Radioactive samples were manufactured by collaboration of three Russian nuclear research Institutes - the Leipunski Institute of Physics and Power Engineering at Obninsk, the Bochvar Institute at Moscow (VNIINN), and the “Maiak” Plant at Ozersk (Chelyabinsk region). Samples are periodically tested for hermetical properties, especially before and after irradiation by alpha activity on exterior surface testing.

Table 1: Basic properties of radioactive samples for transmutation studies

Sample	Decay type	Half-life, y	Weight, g	Purity, %
I-129	β^-	15.7×10^6	0.591	85 I-129
			0.121	15 I-127
Np-237	α	2.14×10^6	1.085	~100 Np-237
Pu-238	α	87.7	0.0477	72.92 Pu-238
				16.75 Pu-239
				2.87 Pu-240
				0.35 Pu-241
				0.11 Pu-242
Pu-239	α	2.41×10^4	0.455	~100 Pu-239

Activation detectors, solid state nuclear track detectors and samples of various technical materials

To determine the neutron field at the places where transmutation samples were located, activation threshold detectors as Al, Co, Cu, Y, Bi, Au where placed on the second section of the blanket.

SSNTDs were used for beam monitoring, investigation of high-energy neutron field ($E > 30$ MeV) between blanket sections, determination of fission abundance and energy output of the blanket, thermal, epithermal and fast neutrons in the produced neutron field, etc.

The investigations of technical properties of superconductor's materials, Hf, Zr, and epoxy (those are of a high importance for accelerator, reactor, and coupled engineering) have also performed.

System of He-3 counters

The basic characteristics of He-3 proportional counter are summarized in Table 2. The measurement system, presented in Fig. 3, consists of a high voltage power supply, a preamplifier suitable for proportional counters (Canberra model 2006), an amplifier (Tennelec model TC205), and a computer based multichannel analyzer (Tennelec PCA III). He-3 was manufacture by LND INC., New York, USA. The system was calibrated using neutrons produced by the Tandem, Van de Graff accelerator facility at the Institute of Nuclear Physics, NCSR Demokritos

(Athens, Greece) [26]. The detector was irradiated with mono-energetic neutrons in the energy range of 230 keV – 7.7 MeV, produced via ${}^7\text{Li}(n,p){}^7\text{Be}$ and ${}^2\text{H}(d,n){}^3\text{He}$ reactions. Due to the high pressure and its large dimensions the He-3 counter could be used effectively for measuring neutron energies up to about 7 MeV. A linear response with incident neutron energy was observed for neutron energies up to this energy, both for the full energy peak and the recoil peak. The resolution varied from 11% for thermal neutrons up to 4% for larger energies.

Table 2. Basic characteristics of the counter.

Detector	Press, atm	Gas content, %	Cathode material Thickness	Anode material diameter	Effective length, cm	Effective diameter, cm
He-3	6	${}^3\text{He}$ 64.7 Kr 33.3 CO ₂ 2.0	Stainless Steel 304 / 0.089 cm	Tungsten / 0.025 mm	15	5

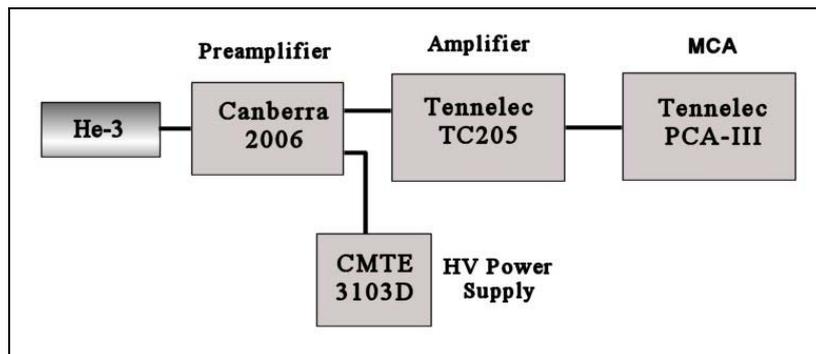


Fig. 3. Neutron counting system (see details in text).

The disadvantage of He-3 counters when they are used in high intensity neutron fields is the relatively high dead time they present, several tenths of μs . In order to avoid space charge effects or even paralyzation of the detector, the maximum count rate should be kept well below 10^4 cps. For this irradiation the motorized stage, which was specifically designed for holding and moving the counter during the experiment, was positioned at the maximum available distance, about 4.7 m from the center of U-blanket (fig. 4).

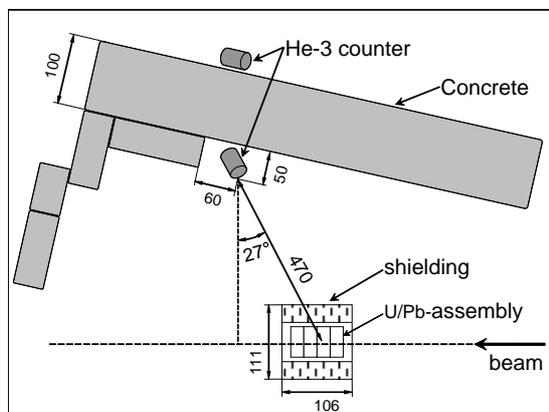


Fig. 4. Arrangement of the counter in respect to the beam direction and U/Pb-assembly of the setup "Energy plus Transmutation" (the distances are in cm).

The cylindrical side of the counter was covered with 1.2 mm Cd to minimize the contribution of scattered thermal neutrons, coming mainly from the concrete walls.

Gamma spectra measurement; installation description

Measurement of activation threshold detectors, Al and Cu beam monitors, and transmutation samples were performed on HPGe detectors provided by Dzhelepov Laboratory of Nuclear Problems of JINR. Description of the main parameters of these detectors is given in table 3. Various geometry position as well as various filters of Pb, Cu, and Cd, were used depending on samples activities. Spectra measurements started few hours after the end of the irradiation and lasted for two weeks (depended from sample to samples). The HPGe detector systems were calibrated using well-defined ^{152}Eu , ^{154}Eu , ^{57}Co , ^{60}Co , ^{137}Cs , ^{88}Y , ^{228}Th radioactive sources. ^{133}Ba source was also used for calibration in several gamma lines ranging from 80 keV up to 2600 keV. The obtained gamma spectra were analyzed and the net peak areas were calculated using the DEIMOS program [27]. All necessary corrections on possible coincidences and background contributions were done. Approximately five hundred of gamma-spectra were measured and analysed.

Results and discussion

This chapter gives some preliminary results of measurements with the ^3He neutron counters, activation threshold detectors from ^{27}Al , ^{87}Y , ^{197}Au , natural U foils, and from SSNT-detectors. Also the transmutation yields results of radioactive nuclear waste isotopes incineration are presented.

(a) Methodical tests of neutron measurements using He-3 counters.

During the first part of the experiment, dedicated to irradiations of emulsions and track detectors, several spectra were collected. In all of them a distortion of the thermal peak (exothermic reaction ${}^3\text{He}(n,p){}^3\text{H}$ has energy $Q = 764$ keV) due to space charge effect is observed, in spite of the relatively smaller intensity of the beam. As an example, the spectra collected during the Polaroid exposure (1 pulse) and during the irradiation for emulsions (6 pulses) are presented in fig. 5. The count rate during these measurements was calculated to be in the range from 14 up to 17 kcps. For dosimetric purposes mainly, during the rest of the irradiation, the counter was placed behind the concrete in a symmetrical position (see Fig 4). The spectrum collected during this irradiation is also presented in Fig 5.

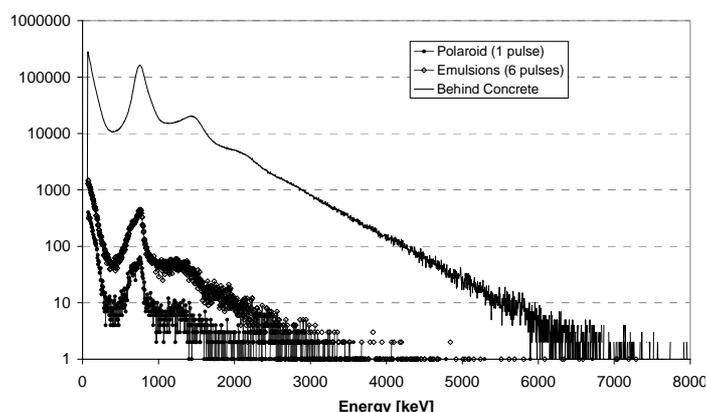


Fig. 5. Spectra collected during the irradiation of “Energy plus Transmutation” setup with deuteron beam at 2.52 GeV

The mean count rate during this measurement was about 20 kcps. The second peak at about 1.5 MeV, present in all spectra, is a summation peak formed when two thermal neutrons are registered simultaneously. According to the above observations, in all the spectra collected the count rate exceeded the limit of this system for neutron spectroscopy. Useful information about neutron spectra of “Energy plus Transmutation” setup could be obtained with smaller intensity and larger duration pulses.

(b) Transmutation of radioactive waste ${}^{129}\text{I}$, ${}^{237}\text{Np}$, ${}^{238}\text{Pu}$ and ${}^{239}\text{Pu}$ in the field of neutrons from ADS. Reactions of radioactive samples – isotopes were irradiated by secondary neutrons generated by the spallation reactions of 2.52 GeV deuterons on lead target. The radioactive samples of ${}^{129}\text{I}$, ${}^{237}\text{Np}$ and stable ${}^{127}\text{I}$ (for detail info see table 1 and fig.1) were irradiated on the top outside surface of the 2nd section of uranium blanket. Transmutation rates of the isotopes – yield of residual nuclei were investigated by gamma-spectroscopy methods. We obtain the data on absolute reaction rate (R-value – number of residual nuclei produced per atom of the sample, per one incident d or p) for some residual nuclei of our samples. ${}^{127}\text{I}$ was used for subtraction of its part (15%) as a contamination of sample ${}^{129}\text{I}$. The results are given in tables 5 and 6 with the data on interaction of

2 GeV protons. Delay between first γ -spectra measurement and end of irradiation (cooling time) was 5 h for protons and 11 h for deuterons. As it is seen from the table 5, the results for deuterons 2.52 GeV (present work) and protons 2.0 GeV [12] are close with small deviations, what means that the main importance was the full energy of the input particles – deuterons and protons.

Table 5. Residual nuclei observed in ^{127}I , ^{129}I , ^{237}Np samples; R-value results for deuteron and proton beam

Residual nuclei	$T_{1/2}$	Deuterons 2.52 GeV (present work)	Protons 2.0 GeV [12]
I-127 sample, R -10 ⁻²⁹			
In-111	16.78 d	0.50(7)	0.38(10)
Te-119	16.03 h	1.15(18)	1.31(27)
Te-119m	4.70 d	1.15(26)	1.03(12)
I-121	2.12 h	3.87(100)	3.13(23)
Sb-122	2.72 d	1.24(15)	-
I-123	13.3 h	11.6(14)	13.0(10)
I-124	4.18 d	18.3(11)	19.0(10)
I-126	13.11 d	70.4(3)	81(4)
I-129 sample, R -10 ⁻²⁹			
Te-121	16.78 d	4.93(94)	-
I-124	4.18 d	4.38(125)	4.0(5)
I-126	13.11 d	10.8(25)	22.5(44)
I-130	12.36 h	816(40)	809(33)
Np-237 sample R -10 ⁻²⁶			
Zr-97	17.0 h	0.188(29)	0.159(8)
Mo-99	2.75 d	1.64(47)	-
Te-132	3.26 d	0.217(32)	0.147(11)
I-133	20.8 h	0.265(75)	0.182(28)
Np-238	2.12 d	17.0(8)	13.3(3)

But for ^{237}Np the yield of residual nuclei for deuterons 2.52 GeV is systematically higher (near factor 1.3) than in case of 2.0 GeV protons.

The experimental values of reaction velocity $R(A)$ and yield of residual nuclei $B(A)$ were calculated by next formulas:

$$R(A) = \frac{N(A)}{n_s \cdot I_d}, \quad B(A) = \frac{N(A)}{m_s \cdot I_d}, \quad R(A) = B(A) \frac{A_s}{N_{Av0}} \quad (2)$$

where $N(A)$ - number of nuclei of isotope A produced in an activation detector, n_s and m_s the number of atoms in the activation detector and its mass in grams, I_d deuteron fluence in the irradiation, N_{Av0} the number of Avogadro. Two important plutonium isotopes (^{238}Pu and ^{239}Pu) were irradiated also on the top of

the second section of the target-blanket system (near the same place where ^{237}Np and ^{129}I and activation threshold samples were placed). Characteristics are given in table 6. As mentioned above, measurements of spectra were started about 11 hours after the end of irradiation. Due to such long cooling time, there is impossible to see produced isotopes with short lifetimes, which were observed in experiments with protons. It is obvious in the case of ^{238}Pu target, in which only two products were found.

Table 6. Residual nuclei observed in ^{238}Pu and ^{239}Pu samples; R- and B-values results for deuteron beam 2.52 GeV

Residual nuclei	$T_{1/2}$	$B (\Delta B) \times 10^5$	$R (\Delta R) \times 10^{27}$
Pu-238 sample			
Zr-97	16.9 h	4.51(11)	15.6(4)
Xe-135	9.14 h	8.0(9)	29(4)
Pu-239 sample			
Ru-103	3.93 d	5.0(4)	19.8(17)
Sb-128	9.01 h	0.18(5)	0.72(22)
Te-132	3.2 d	4.3(4)	16.9(17)
I-133	20.8 h	6.8(7)	27(3)
I-135	6.57 h	4.6(8)	18(3)
Xe-135	9.14 h	3.0(8)	12(3)
Ba-140	12.75 d	5.2(6)	20.4(23)
Ce-143	33.04 h	3.6(4)	14.3(15)
Sr-91	9.63 h	2.6(4)	10.3(17)
Zr-97	16.9 h	5.3(4)	20.9(17)

(c) Experimental studies of spatial neutron distributions using activation detectors. We studied the spatial and energy distributions of neutrons produced at different places of our setup. Produced neutrons induced in activation detectors (foils) γ -decaying products of (n, xn)-, (n, α)- and (n, γ)-reactions. Different thresholds of these reactions allow us to probe energy spectra of neutrons. Foils Al, Y and Au were placed at the different positions of the used setup (also inside the lead target). The activation detectors of the first set were placed at the distances of 0, 11.8, 24.0, 36.2, and 48.8 cm from the front of the Pb target and at the distances from 0.0 to 13.5 cm from the target axis. Measured activities at the end of the bombardment were converted into production rates $B(A)$ of these nuclei (see Eq. 2). Production rate $B(A)$ is very sensitive to the threshold of the reaction and hence also to the neutron spectrum in the foil position. Gamma radiation of each foil was measured two times at different times after irradiation to distinguish the isotopes with different decay times. The results from the analysis of several gamma line intensities from two spectra were used to calculate the experimental production rate B . Weighted averages over the number of spectra

were determined for each individual isotope and foil. Radial distributions of production rates B for several isotopes produced in Au- and Al-foil are shown in the Fig. 6-8 as an example.

Fig. 6. Radial distributions of production rates B for isotope ^{198}Au produced on Au foils by non-threshold (n, γ) reaction for different distances from the front of the Pb-target. The lines are drawn to guide eyes.

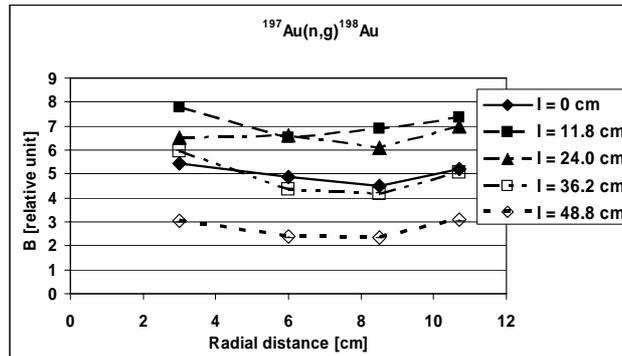
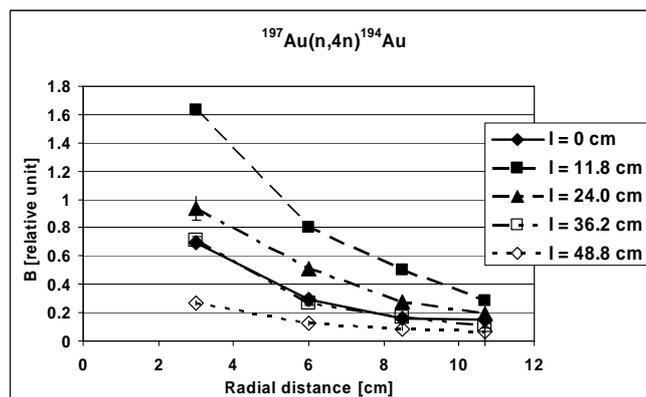
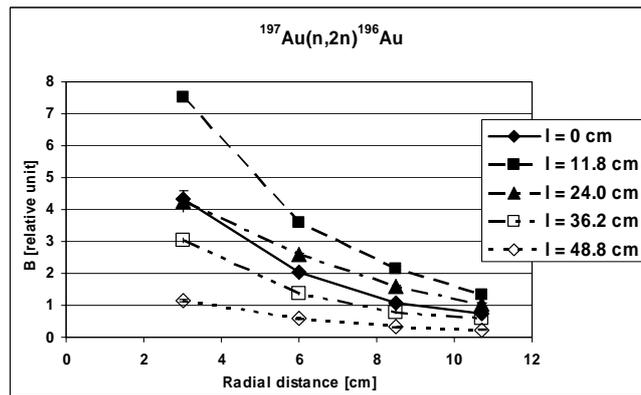


Fig. 7. Radial distributions of production rates B for ^{196}Au (up) and ^{194}Au (down) isotopes produced on Au foils by $(n, 2n)$ and $(n, 4n)$ -reactions with threshold energies 8.1 and 23.2 MeV for different distances from the front of the Pb-target. The lines are drawn to guide the eye.



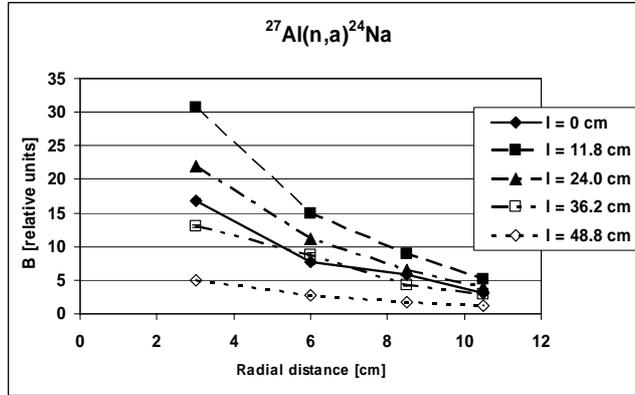


Fig.8. Radial distributions of production rates B for isotope ^{24}Na produced on Al foils by (n,α) -reaction with threshold energy 5.5 MeV. Distributions are shown for different distances l from the front of the Pb-target. The lines are drawn to guide eyes.

General features of observed distributions are as follows:

- they have clearly defined maximum around the distance $l = (16 \pm 2)$ cm from the front of the target;
- the B values are decreasing with the increasing radial distance from the target axis (beam axis position);
- while B value distributions for non-threshold (n, γ) -reaction are almost flat, corresponding distributions in the case of threshold reactions changed over magnitude and in case of radial distributions are close to exponential form.

Therefore, one can conclude that epithermal and resonance neutron fluences are almost constant in the target volume and at the same time high energy neutron fluence depend strongly on the radial distance from the target axis. These features are very similar to our previous experimental results obtained using proton beams [5,28,29] and they fit together with our systematic of proton experiments.

The comparison between experimental and simulated production rates of many threshold reactions makes possible to test accuracy of the description of neutron production in the wide neutron energy interval. It is also possible to restore the neutron energy spectrum using cross sections dependence on neutron energy for sufficient number of the production reactions.

(d) Measurements of neutron spatial-energy distribution using SSNTDs.

The spatial-energy distribution of neutrons was also measured by SSNTDs. This part includes measurement of distribution of fission rates of ^{238}U

and ^{238}U , reactions of radiation capture of neutrons by ^{238}U , and also a spectral index $\frac{\sigma_{\text{capt}}^{238\text{U}}}{\sigma_f^{238\text{U}}}$ on radius of the U/Pb-assembly. Knowing the yield of fission products [30], it is possible to determine distribution of fission

density of ^{238}U . It is interesting to compare results received by two independent experimental methods among themselves and also with calculation results. It is obvious from Fig. 9a, that results of measurement of radial distribution of number of fission reactions ^{238}U determined by two experimental methods are in coincidence in the limit of experimental errors already since distance of 30 mm from longitudinal Pb target axis. Calculation also well describes fission process in a blanket material (natural uranium) and on periphery of the assembly.

The spectral index characterizes a ratio between speeds of capture and fission of neutrons in a material uranium blanket. Results of experiment and calculation of a spectral index coincide in limits of experimental error (Fig. 10). The difference (1.5 times) is observed on periphery of the assembly.

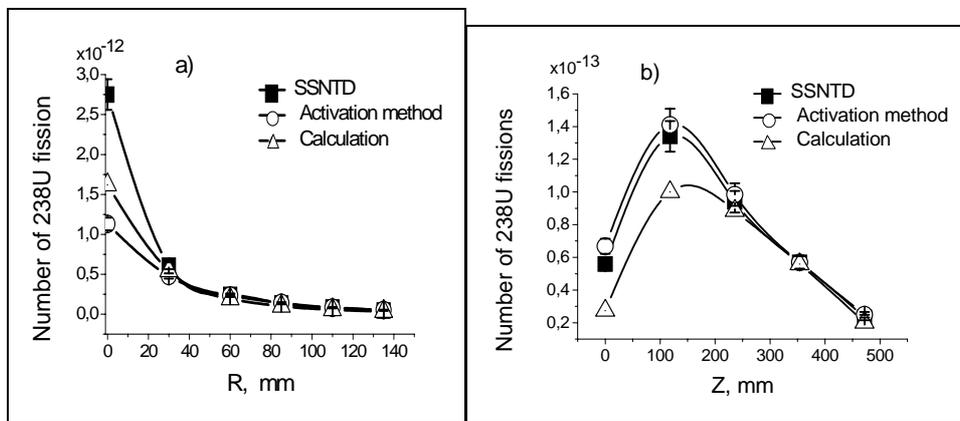
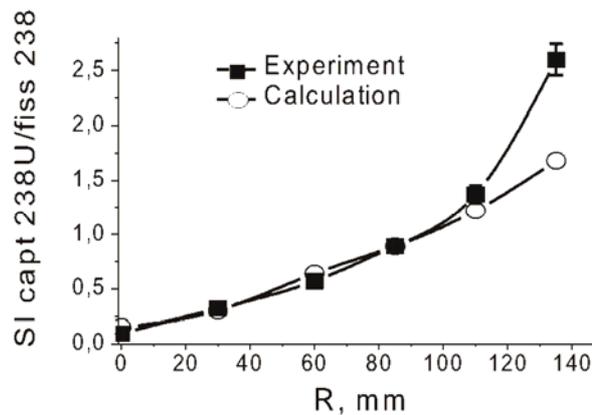


Fig. 9. Radial (distance $Z = 118$ mm) and axial (distances $R=85$ mm from the target axis) distributions of number of fission reactions ^{238}U . The data were normalized on one nucleus ^{238}U and one incident deuteron.

Fig. 10. Radial distribution of a spectral index (distance $Z = 118$ mm from the target front). Lines connecting the data points are drawn to guide the eye.



As it was marked above, it speaks underestimation at calculations of influence of the neutrons moderated in biological shielding and reflected by biological shielding. It is obvious from Fig.10 that the number of ^{238}U fissions exceeds number of radiation captures of neutrons in three times on the border of a lead target and blanket ($R=32$ mm). Processes of radiation capture of neutrons begin to prevail with increasing radial distance, because in the process of moderation of neutrons as a result of not elastic collisions with nuclei of a material blanket. The developed combined track with γ -spectrometry technique of the spectral index determination provides reception of the information from the same sample by SSNTD methods (fission tracks density of ^{238}U) and by a γ -spectrometer method (on a γ -line nuclide ^{239}Np with energy 277.6 keV) and which allows to measure spectral index with the error no more than 15 %. The developed technique will allow determination of ^{239}Pu accumulation in the U-blanket. Experimental value of total mass of ^{239}Pu accumulated in the setup is $1.6(2)\cdot 10^{-8}$ g. We obtained also the value $1.43\cdot 10^{-8}$ g using MCNPX code [31]. Calculation shows good agreement with experimental result within the limits of errors.

Fast neutron distribution along the U-blanket. The neutron flux on the U-blanket surface reflects neutron production along the Pb-target from spallation reactions by primary and secondary particles and also neutron production in the uranium blanket by secondary particles. According to the experimental data, thermal neutrons were not detected because no difference between the tracks density on the Cd-covered and the Cd-uncovered regions of CR39 with $\text{Li}_2\text{B}_4\text{O}_7$ converter was observed. The same conclusion was verified by the supplementary measurements of thermal neutrons using ^{235}U fission detectors [6]. The track density of the Cd-covered region of CR39 with $\text{Li}_2\text{B}_4\text{O}_7$ converter corresponds to epithermal neutron fluence (up to 10 keV) which presented to be one order of magnitude less than fast neutron fluence (between 0.3-3 MeV) detected by proton recoil on the CR39 detector itself. Thus, it is obvious that the U/Pb-assembly produces mostly intermediate-fast neutrons.

Fission detectors with ^{232}Th converter were also used in order to measure fast neutrons fluence above 2 MeV. The fast neutron distributions along the U-blanket measured by both methods are in good agreement as presented in Fig.11. In addition, approximately the same values of the neutron fluxes have been calculated using the high-energy transport code DCM-DEM (Dubna). The fast neutron production increases up to about 15 cm from the beam entrance and then decreases along the target as presented in Fig.12. The same intensity distribution and approximately the same intensities of fast neutrons have been determined in an irradiation of the same spallation source by a 1.5 GeV proton beam [32].

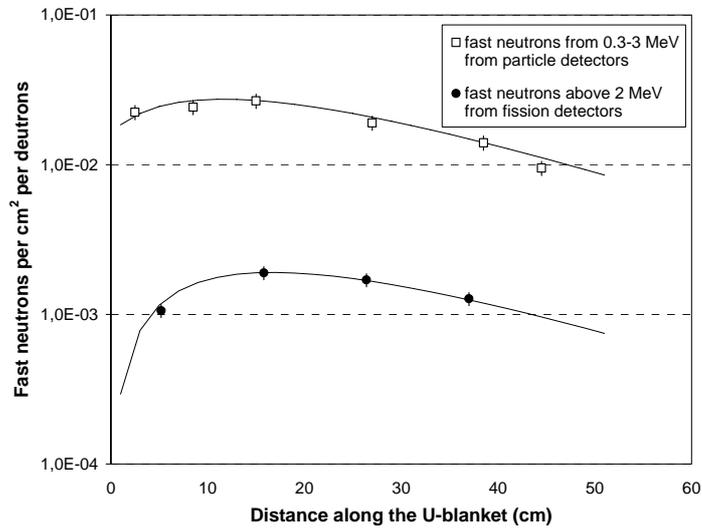


Fig.11. The fast neutron distribution along the U-blanket measured by particles and fission detectors.

The fluence of thermal and fast neutrons escaping the shielding also was measured with CR39 detectors. The intensity of fast neutrons escaping the polyethylene was two orders of magnitude less than on the U-blanket surface (Fig 12). While less than half of these neutrons are in the thermal energy range about 2×10^{-5} neutrons per cm^2 per deuterons. Therefore the polyethylene shielding proves to be an efficient moderator thermalizing a large amount of the fast neutrons produced by the spallation neutron source.

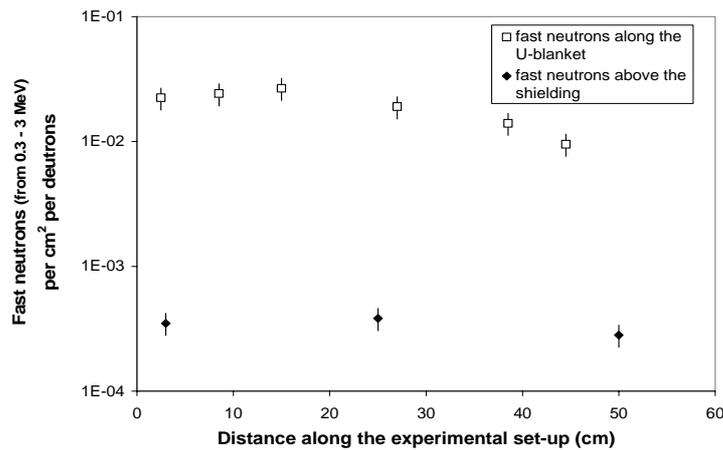


Fig. 12. Neutron distribution escaping the polyethylene shielding and comparison with neutron distribution along the U-blanket surface.

Conclusion

We performed the first experiment on investigations of transmutation (incineration) of radioactive waste atomic reactors (^{129}I , ^{237}Np , ^{238}Pu and ^{239}Pu) using deuterons beam with 2.52 GeV energy. Methodical measurements with the He-3 counters were done to determine the neutron fluences. Neutron spatial distribution was studied using activation threshold detectors. Natural uranium foils were used to determine the (n,γ) -reaction rate along the target length and radius. Spectral indexes (capture to fission ratio) were also determined. Experimental results were compared with MCNPX simulations. Experimental and simulated values of spectral indexes agree well. Total experimental value of the ^{239}Pu production in the whole blanket was estimated as $1.6(2)\cdot 10^{-8}$ g, the theoretical prediction using MCNPX is $1.43\cdot 10^{-8}$ g.

The measured experimental data complement our systematic obtained using proton beams with energies from 0.7 to 2.0 GeV.

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